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Infinitesimal Transformations in the Cartan Space

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矢野 晋平 : Cartan 空間の無限小變換

Introduction

The geometrical property of continuous groups of transformations in the Riemannian space, is one of the most interesting in the theory of groups¹⁾. K.YANO²⁾ and E.T.DAVIES³⁾ extended this property by making use of the *Lie* derivatives in to generalized spaces. On the other hand, in investigation of the transformations in the *Finsler* space, M.S.KNEBELMAN⁴⁾ and E.T.DAVIES⁵⁾ formulated many interesting theories which E.T.DAVIES⁶⁾ established in the metric space.

In the present paper we show whether analogous theories exist or not in the *Cartan* space.⁷⁾

§ 1. Fundamental equations in the n-dimensional Cartan space.

In an n -dimensional manifold, if any $(n-1)$ -dimensional sub-space is given by the equations

$$(1.1) \quad x^i = x^i(y^1, \dots, y^{n-1}) \quad (i = 1, \dots, n),$$

and if the $(n-1)$ -dimensional area of a region R_{n-1} in the sub-space, is given by the $(n-1)$ -ple integral

$$(1.2) \quad O = \int_{n-1} F(x, \partial x / \partial y) dy^1, \dots, dy^{n-1},$$

then the n -dimensional manifold is called the *Cartan* space⁸⁾.

In this space we define the elements of hypersurface u_α by the following determinants made of (1.1)

$$(1.3) \quad u_\alpha = (-1)^{\alpha+1} \begin{vmatrix} \frac{\partial x^1}{\partial y^1} & \dots & \frac{\partial x^{\alpha-1}}{\partial y^1} & \frac{\partial x^{\alpha+1}}{\partial y^1} & \dots & \frac{\partial x^n}{\partial y^1} \\ \vdots & & \vdots & \vdots & & \vdots \\ \frac{\partial x^1}{\partial y^{n-1}} & \dots & \frac{\partial x^{\alpha-1}}{\partial y^{n-1}} & \frac{\partial x^{\alpha+1}}{\partial y^{n-1}} & \dots & \frac{\partial x^n}{\partial y^{n-1}} \end{vmatrix}.$$

Let $T^i_j(x, u)$ be the components of a mixed tensor of degree zero in the u , then the two kinds of differential operators⁹⁾ $|_k$ and $|^k$ are applied to the T^i_j in the following expressions

- 1) See, [7]
- 2) See, [10], §11
- 3) See, [4]
- 4) See, [8]
- 5) See, [5]
- 6) See, [6]
- 7) The notations in our paper, are the same as those used by L.BERWALD [1], and K.YANO [9], omitting * in T^i_j .
- 8) See, [1] Einleitung, see [2] §6
- 9) See, [1] §8, [9] §8

$$(1.4) \quad T^i_j |^k = T^i_{j,k} + T^i_j ||^\beta \Gamma_{\beta 0 k} + T^\beta_j \Gamma_{\beta}^i k - T^i_\beta \Gamma_j^{\beta k},$$

$$(1.5) \quad T^i_j |^k = T^i_j ||^k + T^\beta_j A_{\beta}^{ik} - T^i_\beta A_j^{\beta k},$$

where

$$(1.6) \quad l_i = (\sqrt{g}|L)u_i,$$

$$(1.7) \quad T^i_j |^{k,h} = (L|\sqrt{g})T^i_j |^{k,h},$$

Throughout the paper we shall use the notations, $T^i_{j,k} = \partial T^i_j / \partial X^k$, $T^i_j |^{k,h} = \partial T^i_j |^{k,h} / \partial u_k$.

In the Cartan space there are four kinds of curvature tensors¹⁾ given by the following expressions

$$(1.8) \quad K_{i^j k h} = 2(\Gamma_{i^j [k, h]} + \Gamma_{i^j [k} \Gamma_{h]}^j \alpha + \Gamma_{i^j [k} ||^\alpha \Gamma_{h]}^j \alpha),$$

$$(1.9) \quad R_{i^j k h} = K_{i^j k h} - A_{i^j}^\alpha K_{\alpha 0 k h},$$

$$(1.10) \quad P_{i^j k}^{h, h} = \Gamma_{i^j k} ||^h - A_{i^j}^{h, h} - A_{i^j}^\alpha l_\beta \Gamma_{\alpha}^{\beta k} ||^h,$$

$$(1.11) \quad S_i^{j k h} = 2 A^{\alpha j [k} A_{\alpha} i^{h]}.$$

§2 The relations of the differential operators to the curvature tensors.

On consideration of (1.7), from (1.4), (1.5), (1.8) and (1.9) we have

$$(2.1) \quad 2 T^i_j |^{[k h]} = T^i_j ||^\alpha K_{\alpha 0 h k} + T^{\alpha j} K_{\alpha}^i h k - T^i_\alpha K_j^{\alpha h k} = T^i_j ||^\alpha R_{\alpha 0 h k} + T^{\alpha j} R_{\alpha}^i h k - T^i_\alpha R_j^{\alpha h k}.$$

By virtue of the relation²⁾ $(L|\sqrt{g}) ||^i = L|\sqrt{g} (l^i - A^i)$ we see that

$$(2.2) \quad T^i_j |^{[k h]} = T^i_j ||^{[k} (l^{h]} - A^{h]}).$$

Put the metric tensors g_{ij} and g^{ij} in to these equations and in consequence of the relations³⁾ $g^{ij} ||^h = -2A^{ijh}$ and $g_{ij} ||^h = 2A_{ij}^h$ we have

$$(2.3) \quad A_i^{j[k} ||^{h]} = 2A^{\alpha j [k} A_{i \alpha}^{h]} + A_i^{j[k} (l^{h]} - A^{h]}).$$

Since $\Gamma_j^i k ||^h$ is a tensor of order four⁴⁾, from the connection parameters $\Gamma_j^i k$ we have

$$(2.4) \quad \Gamma_j^i m ||^{[k h]} = \Gamma_j^i m ||^{[k} (l^{h]} - A^{h]}).$$

If we substitute the right side of (1.6) in the relations⁵⁾

$$(2.5) \quad l_{i,j} + l_i ||^\alpha \Gamma_{\alpha 0 i} - \Gamma_{i 0 j} = l_{i j} = 0,$$

in use of the equations⁶⁾

$$(2.6) \quad l_i ||^\alpha = \delta^\alpha_i - l_i (l^\alpha - A^\alpha), \quad l_i ||^\alpha = g^{i\alpha} - l^i (l^\alpha + A^\alpha),$$

we have $(L|\sqrt{g})_{,h} + (l^\alpha - A^\alpha) \Gamma_{\alpha}^{\beta h} u_\beta = 0$. From these equations and (2.2) we obtain finally

$$(2.7) \quad T^i_j |^{[k} ||^h - T^i_j ||^{[k} ||^h = T^i_j ||^\alpha \Gamma_{\alpha}^{\beta h} ||^k + T^{\alpha j} \Gamma_{\alpha}^i h ||^k - T^i_\alpha \Gamma_j^{\alpha h} ||^k.$$

Since the necessary and sufficient condition for that a Cartan space become an affine connected space is that $\Gamma_{\alpha}^i j ||^k = 0$ ⁷⁾, from (2.7) we get

Theorem [2.1]. If the Cartan space is an affine connected space, then two operators $|^h$ and $||^k$ may be interchangeable.

By virtue of (2.2) and (1.5) we have

$$(2.8) \quad T^i_j |^{k,h} - T^i_j ||^{h,k} = 2 T^i_j ||^{[k} (l^{h]} - A^{h]} - T^i_j ||^\alpha A_{\alpha}^{h k} + T^{\alpha j} A_{\alpha}^{i k} ||^h - T^i_\alpha A_j^{\alpha k} ||^h.$$

1) See [1], §12 and [9], §9.
 2) See [1], §5
 3) See [1], §5
 4) See [1], §16
 5) See [1], §8
 6) See [1], §5
 7) See [9], §12

Since $T^i_{j|h}$ and $T^i_j{}^{|k}$ are the tensors, from (2.7) and (1.10) we have

$$(2.9) \quad T^i_{j|h}{}^{|k} - T^i_j{}^{|k}{}_h = T^i_j{}^{|\alpha} T_{\alpha}{}^{\beta}{}_{|h} l^k_{\beta} - T^i_j{}^{|\alpha} A_h{}^{\alpha k} + T^{\alpha}{}_j P_{\alpha}{}^k{}_h - T^i_{\alpha} P_j{}^{\alpha k}.$$

Considering (1.5), (2.2), (2.3) and (1.11) we find

$$(2.10) \quad 2T^i_j{}^{[k|h]} = 2T^i_j{}^{[k}(l^{h]} - A^{h])} + T^{\alpha}{}_j S_{\alpha}{}^{ikh} - T^i_{\alpha} S_j{}^{\alpha kh} + 2T^{\alpha}{}_j{}^{[k} A_{\alpha}{}^{]hk]} - 2T^i_{\alpha}{}^{[k} A_j{}^{\alpha]hk}.$$

From these equations we have

$$(2.11) \quad 2T^i_j{}^{[k|h]} = 2T^i_j{}^{[k}(l^{h]} - A^{h])} + 2T^i_j{}^{|\alpha} A_{\alpha}{}^{[k|h]} + T^{\alpha}{}_j S_{\alpha}{}^{ikh} - T^i_{\alpha} S_j{}^{\alpha kh}.$$

Taking l_i in place of T^i_j in the above equations and considering the relations $S_{i_0}{}^{kh} = 0^{(D)}$ and $l_i{}^j = (\delta^j_i - l^j l^i)^{(D)}$, we find $A_i{}^{[jk]} = \delta^{[j} l^k] A^i$. From these equations we have

Theorem [2.2]. If the condition $A^i = 0$ is satisfied, then the tensor A^{ijk} is symmetric in the indices i, j, k .

§ 3 The Lie derivatives

Let us consider a one-parameter continuous group of transformations defined by finite equations $\bar{x}^i = \bar{x}^i(x, \alpha)$, and the infinitesimal transformations of the group given by

$$(3.1) \quad \bar{x}^i = x^i + \xi^i(x) \delta t,$$

where δt is an infinitesimal constant and ξ^i a contravariant vector field depending on a point.

If we introduce the notations $\partial x^i / \partial y^{\alpha} = x^i_{\alpha}$ and $\partial x^i_{\alpha} = \bar{x}^i_{\alpha} - x^i_{\alpha} = \xi^i_{,j} x^j_{\alpha} \delta t$, from (1.3) we have

$$(3.2) \quad \sum_{\alpha=1}^{n-1} (\partial u_i / \partial x^j_{\alpha}) x^k_{\alpha} = u_i \delta^k_j - u_j \delta^k_i.$$

Put the elements of hypersurface deformed by the transformation (3.1) in the form $\bar{u}_i = u_i + \delta u_i$, where $\delta u_i = (\partial u_i / \partial x^j_{\alpha}) \delta x^j_{\alpha}$. On neglecting terms of higher orders, considering (3.2) we have

$$(3.3) \quad \delta u_i = (u_i \delta^k_j - u_j \delta^k_i) \xi^j_{,k} \delta t.$$

If we define, with reference to the transformation (3.1), the *Lie* derivative of any geometric object T : (e.g. a tensor, tensor density, connexion parameters, etc.) as

$$X T :: = \lim_{\delta t \rightarrow 0} \frac{T :: (\bar{x}, \bar{u}) - T :: (x, u)}{\delta t}.$$

In consequence of the relations $\partial \bar{x}^i / \partial x^j = \delta^i_j + \xi^i_{,j} \delta t$ and $\partial x^j / \partial \bar{x}^i = \delta^j_i - \xi^j_{,i} \delta t$, from (3.3) the Lie derivative of tensor T^i_j can be written in the form¹⁾

$$(3.4) \quad X T^i_j = T^i_{j,\alpha} \xi^{\alpha} - T^i_j{}^{|\alpha} \varphi_{|\alpha} + T^i_{\alpha} \xi^{\alpha}_{,j} - T^{\alpha}{}_j \xi^i_{,\alpha}, \text{ where } \varphi = \xi^i l_i.$$

If $T^i_j(x, u)$ are the components of a tensor density of weight p , considering the equations $|\partial x / \partial \bar{x}|^p = 1 - p \xi^{\alpha}_{,\alpha} \delta t$, we have

$$(3.5) \quad X T^i_j = T^i_{j,\alpha} \xi^{\alpha} - T^i_j{}^{|\alpha} \varphi_{|\alpha} + p T^i_j{}^{\beta}{}_{|\beta} \xi^{\alpha} - T^{\alpha}{}_j \xi^i_{,\alpha} + T^i_{\alpha} \xi^{\alpha}_{,j}.$$

Let \mathfrak{S}^i_j be the components of a tensor of weight p and homogeneous of degree q in the u , then we have

$$(3.6) \quad X \mathfrak{S}^i_j = \mathfrak{S}^i_{j,\alpha} \xi^{\alpha} - \mathfrak{S}^i_j{}^{|\alpha} \xi^{\beta}_{,\alpha} l_{\beta} + \mathfrak{S}^{\alpha}{}_j \xi^i_{,\alpha} + \mathfrak{S}^i_{\alpha} \xi^{\alpha}_{,j} + (p+q) \mathfrak{S}^i_j{}^{\alpha} \xi^{\alpha}_{,\alpha}.$$

1) See [1] §12
 2) See [9] §8
 3) See [6] §2
 4) See [6] §2

Since u_i is a covariant vector density of weight -1 and dx^i a contravariant vector. From (3.4) and (3.5) we get $Xu_i=0, Xdx^i=0$, consequently we have

Theorem [3.1]. The Lie derivatives of the hypersurface and line elements are equal to zero.
Similarly, we have

Theorem [3.2]. The Lie derivatives of ξ^i and δ^i_j are equal to zero.

If we substitute l_i in (3.4) and considering (2.5) and (2.6) we have $Xl_i=l_i(l^\alpha-A^\alpha)\theta_{i\alpha}$.

If these equations be contracted by l^i , we have

$$(3.7) \quad Xl_i=l l^\alpha Xl_\alpha.$$

Substitute the right side of (1.6) in (3.7), from theorem [3.1] we have

$$(3.8) \quad (L/\sqrt{g})X(\sqrt{g}/L)=l^\alpha Xl_\alpha.$$

From the parameters Γ^i_{jk} , considering the equations $\xi^i_{,j}=\xi^i_{,j}+\Gamma^\alpha_{j\alpha}\xi^i$, we obtain

$$(3.9) \quad X\Gamma^i_{jk}=\xi^i_{,jk}+\xi^\alpha_{,j}K^\alpha_{jk}-\Gamma^i_{k\alpha}\xi^\alpha_{,j}+\Gamma^i_{j\alpha}\xi^\alpha_{,k}.$$

Since Γ^i_{jk} is a tensor, we have

Theorem [3.3]. The Lie derivative of connection parameter Γ^i_{jk} is a tensor.

The application to the metric tensor g_{ij} will therefore give

$$(3.10) \quad Xg_{ij}=\xi_{i,j}+\xi_{j,i}-2A_{ij}^\alpha\theta_{i\alpha}.$$

The equations of Killing, in the Riemannian geometry, are given by the relation

$$(3.11) \quad Xg_{ij}=0.$$

Therefore, we may say the transformation (3.1) giving (3.11) an infinitesimal motion.

For the motion, considering (3.11), (3.7) and the relation $X(l^i l_i)=0$, we have

$$(3.12) \quad Xl_i=0.$$

§ 4 Alternation of Lie derivative and the differential operators.

The following alternate formula will be needed later.

From (2.1), (2.7) and (3.9) we obtain

$$(4.1) \quad X(T^i_{jk})-(XT^i_j)_k=T^i_{j\alpha}\xi^\alpha_{,k}+T^\alpha_{jk}X\Gamma^\alpha_{ik}-T^i_{\alpha k}X\Gamma^\alpha_{j\alpha}.$$

These equations imply

Theorem [4.1]. If infinitesimal transformation (3.1) satisfy conditions $X\Gamma^i_{jk}=0$, then the Lie derivative and the operator $|_k$ are interchangeable.

If \mathfrak{S}^i_j is a tensor density of weight p and homogeneous of degree q in the u , then $\mathfrak{S}^i_{j;k}$ is a tensor of weight $p+1$ and degree $q-1$. Therefore considering (3.6) we have

$$(4.2) \quad X(\mathfrak{S}^i_{j;k})-(X\mathfrak{S}^i_j)_k=0.$$

Theorem [4.2]. Let \mathfrak{S}^i_j be a tensor of any weight and degree in the u , then the Lie derivative and partial differentiation with u are mutually interchangeable.

From (3.8) and (4.2) we have

$$(4.3) \quad X(\mathfrak{S}^i_{j;k})-(X\mathfrak{S}^i_j)_k=-\mathfrak{S}^i_{j\alpha}l^\alpha Xl_k.$$

from this equation and (1.7) we have

Theorem [4.3]. If \mathfrak{S}^i_j is independent of u , then $X\mathfrak{S}^i_j$ is same also.

1) See [6], §2

By virtue of (3.12) we obtain

Theorem [4.4]. If (3.1) is a motion, the Lie derivative and the operator $\|^{k}$ are interchangeable.

Finally we have from (4.3)

$$(4.4) \quad X(T^l_j{}^k) - (XT^l_j)\|^{k} = -T^l_j\|^{k}l^\alpha Xl_\alpha + T^\alpha_j XA_\alpha{}^{kk} - T^l_\alpha XA_j{}^{\alpha k}.$$

Since $\Gamma_j{}^k{}_{;h}$ is a tensor density of weight 1 and of the homogeneous degree -1 in the u , then considering (4.2) and (4.3), we have

$$(4.5) \quad X(\Gamma_j{}^k\|^{h}) - (X\Gamma_j{}^k)\|^{h} = -\Gamma_j{}^k\|^{h}l^\alpha Xl_\alpha.$$

From the relation $\Gamma_j{}^k{}_{;h} = \Gamma_j{}^k + X\Gamma_j{}^k\delta l$, $K_j{}^k{}_{;h} = 2\Gamma_j{}^k{}_{[c;h]} + 2\Gamma_j{}^{\alpha}{}_{[c;h]}l_\alpha + 2\Gamma_j{}^k{}_{[c;h]}\|^{a}T_{h]0\alpha}$, considering (4.5) and (3.7), we find

$$(4.6) \quad XK_j{}^k{}_{;h} = 2(X\Gamma_j{}^k{}_{[c;h]}) + 2\Gamma_j{}^k{}_{[c;h]}\|^{a}X\Gamma_{h]{}^\beta}{}_{\alpha}l_\beta.$$

From these equations we have

Theorem [4.5]. If the transformations (3.1) satisfy the relations $X\Gamma_j{}^k = 0$, the Lie derivative of $K_j{}^k{}_{;h}$ vanishes.

Substitute g_{ij} in (4.3), we have

$$(4.7) \quad 2XA_{ij}{}^v = -2A_{ij}{}^k l^\alpha Xl_\alpha + (Xg_{ij})\|^{k},$$

from these equations and (3.12) we have

Theorem [4.6]. When the space admits an infinitesimal motion (3.1), the Lie derivative of connexion $A_{ij}{}^k$ vanishes.

Put g_{ij} in (4.1), from the relation $g_{ij}{}^k = 0$ we have

$$(4.8) \quad (Xg_{ij})_{;k} = (g_{\beta j}\delta^{\alpha}{}_{\beta} + g_{i\beta}\delta^{\alpha}{}_{\beta} - 2A_{ij}{}^{\alpha}l_\beta)X\Gamma_{\alpha}{}^{\beta}{}_{;k}{}^{1)}$$

From these equations we have

Theorem [4.7]. When the space admits an infinitesimal motion (3.1), the Lie derivative of connexion $\Gamma_j{}^k$ is equal to zero.

§ 5 Groups of motions and affine motions

In this chapter, we shall study the continuous groups of motions, then from theorems [4.6], [4.7] and (3.12) we have

$$(5.1) \quad XA_i{}^{jk} = 0, \quad X\Gamma_i{}^j{}_{;k} = 0, \quad Xl_i = 0.$$

By virtue of (3.11) and theorem [3.1] we see that $X(ds) = X(g_{ij}dx^i dx^j) = 0$, hence we have

Theorem [5.1]. When the space undergoes a motion into itself, lengths and angles remain unaltered.

A necessary and sufficient condition that the trajectories of two motions be the same is that (3.9) be satisfied by $\bar{\xi}^i$ and ξ^i , where $\bar{\xi}^i = \rho\xi^i$, then we have $\rho_{,\alpha}(2A_{ij}{}^{\alpha}\xi^i l_\beta - \delta^{\alpha}{}_{\beta}\xi^j - \delta^{\alpha}{}_{j}\xi^i) = 0$, if the matrix of the coefficients of $\rho_{,\alpha}$ is of rank n , we obtain $\rho_{,\alpha} = 0$, that is $\rho = \text{constant}$: we have

Theorem [5.2]. Two groups of motions cannot have the same trajectories.

If the space admits absolute parallelism, then denoting by $h^i{}_\alpha$ the n contravariant vectors of an absolutely parallel ennuple,²⁾ the suffixes i and j being a component and α and β being a number of the vectors, we have

$$(5.2) \quad h^i{}_\alpha h_{\beta i} = \delta_{\alpha\beta}, \quad \sum_{\alpha=1}^n h^i{}_\alpha h_{\alpha j} = \delta^i_j.$$

1) See [6] §2

2) See [10] §3

From the relations $Dh^l_a = h^l_{a,j} dx^j + h^l_{a^j} D_l j = 0$, we have

$$(5.3) \quad h^l_{a^j} = 0, \quad h^l_{a^j} = 0, \quad (i, j, a = 1, \dots, n).$$

Put h^l_a in to (2.1) and considering (5.2) and (5.3), we have $R^j_{k^h} = 0$ and $K^j_{k^h} = 0$. Similarly, from (2.9) and (2.11) we have $P^j_{k^h} = 0$ and $S^j_{k^h} = 0$, consequently we have

Theorem [5.1]. If the space admits absolute parallelism, all the curvature tensors vanish.

From (5.2) and (5.3) we have

$$(5.4) \quad \Gamma^j_{i^k} = \sum_{a=1}^n h^j_a h_{a^i, k} + A_i^{j\beta} \Gamma_{\beta 0^k}, \quad h_{a^i}{}^j = h_{a^k} A_i^{k^j}.$$

Let us multiply l_j in the first equation of (5.4) and contract on j , then $\Gamma_{\beta 0^k}$ can be solved, considering the relations $A_i^{j\beta} l_j = l_i A_\beta$ and $(\delta^{\alpha_j} - l_j A^\alpha) \delta^i_\alpha - l_\alpha A^i = \delta^i_j$,

Putting the obtained value of $\Gamma_{\beta 0^k}$ in to the equation again, we have

$$(5.5) \quad \Gamma^j_{i^k} = \sum_{a=1}^n h^j_a h_{a^i, k} + \sum_{a=1}^n A_i^{j\alpha} h^\beta_{a^i} h_{a\alpha, k}$$

and from the second equation (5.4) we have

$$(5.6) \quad A_i^{j^k} = \sum_{a=1}^n h_{a^i}{}^k h^j_a.$$

Suppose that the space admits an infinitesimal transformation (3.1) satisfying the equations (5.1), from (5.3) and theorem [4.1] we have

$$(5.7) \quad (Xh^l_a)_j = 0.$$

Since Xh^l_a is a vector, there are certain scalars C^b_a satisfying the relations

$$(5.8) \quad Xh^l_a = C^b_a h^l_b.$$

From (5.7) and (5.3) we have $C^b_{a^j} = 0$, this condition may be written in the equation

$$(5.9) \quad C^b_{a^j} + C^b_{a^i} \Gamma_{a 0^j} = 0.$$

Similarly, from (5.1), (5.2) and (4.4) we have $C^b_{a^i} = 0$, that is $C^b_{a^i} = 0$, considering these relations and (5.9), we have

Theorem [5.2]. If the space of absolute parallelism admits an infinitesimal transformations (3.1) satisfying the equations (5.1), the Lie derivatives of the n independent vectors of an absolutely parallel ennuple are expressed by linear combinations of them with constant coefficients.

If the transformation (3.1) satisfies the conditions $X\Gamma^i_{j^k} = 0$, then we may say, following K.YANO¹⁾ and M.S.KNEBELMAN²⁾, the transformation is an *affine motion*.

Let $\xi^i_{(\alpha)}$ and $\xi^i_{(\beta)}$ being the components of two linearly independent vectors giving the infinitesimal affine motions in our space, then the components of their alternations are $\xi^m_{(\beta)} \xi^i_{(\alpha), m} - \xi^m_{(\alpha)} \xi^i_{(\beta), m} = X_{(\beta)} \xi^i_{(\alpha)}$, in our notation $\xi^i = X_\beta \xi^i_{(\alpha)}$.

We wish to prove that, ξ^i are the components of a vector giving an affine motion.

From the conditions $X_{(\alpha)} \Gamma^i_{j^k} = 0$ and $X_{(\beta)} \Gamma^i_{j^k} = 0$, considering theorems [4.1] and [4.5], for ξ^i the equations (3.9) may be written in the form

$$X\Gamma^i_{j^k} = (X_{(\beta)} \Gamma^i_{k^j}) l_\rho + \Gamma^i_{k^j} X_{(\beta)} l_\rho \xi^\rho_{(\alpha), \alpha}.$$

From (4.5) and (3.7), the right side of above equation vanishes, then we have

Theorem [5.3]. If $D_\alpha f$ for $\alpha = 1, 2, \dots, p$ are generators of p one-parameter groups of the affine motions, so also are each of the commutators $(D_\alpha D_\beta) f$, for $\alpha, \beta = 1, 2, \dots, p$.

1) See [10], §3

2) See [8], §2

From this result and the fundamental theorems of continuous groups we have

Theorem [5.4]. If $D_{(a)}f$ for $a=1, \dots, p$ are generators of the complete set of one-parameter groups of affine motions, they are the generators of a group G_p of affine motions.

(March, 1952)

References

- [1] L.BERWALD Über die n -dimensionalen Cartansche Räume und eine Normalform der zweiten Variation eines $(n-l)$ -fachen Oberflächen integrals. Acta Math, 71 (1939).
- [2] E.CARTAN Les espaces de Finsler, Actual. Sci. Ind. 79, (1933).
- [3] E.CARTAN Les espaces métriques fondés sur la notion d'aire, Actual Sci. Ind 72 (1933).
- [4] E.T.DAVIES On the infinitesimal deformations of a space, Annali di Matematica 12 (1933-1934) applicata 12 serie 4.
- [5] E.T.DAVIES Lie derivation in generalized metric spaces, Trtolini annali di Matematica 18 Serie 4 (1939).
- [6] E.T.DAVIES Motions in a metric space based on the notion of area, J.Math. Oxford 16 (1945).
- [7] L.P.EISENHART Continuous groups of transformations, Oxford University press (1933).
- [8] M.S.KNEBELMAN Collineations and motions in generalized space, Amer. J.Math. 51 (1929).
- [9] K.YANO Les espaces de Cartan, Tokoyo, Butsurigakko zassi (1943).
- [10] K.YANO Groupes of transformations in generalized space, Akademie presscompany Tokyo, Japan (1949).
- [11] K.YANO Sur la théorie des déformations infinitésimales. J.Fac. Imp. Univ. Tokyo Sect. I vol VI. 1949 (1-75).