



Connection Parameters in a Special Kawaguchi Space

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Connection Parameters in a Special Kawaguchi Space

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叶 長太郎：特殊河口空間における接続経数について

Introduction

Geometry in a special Kawaguchi space with the arc length,

$$S = \int \{A_i(x, x'')x''^{ii} + B(x, x', x'')\}^{1/2} dt,$$

was studied by T. Ohkubo [1].

In the present paper we shall introduce the other covariant derivative which is invariant for a change of parameter. In this space we have the Synge vectors

$$\overset{2}{E}_i = 3A_{i(1)j} x'^j + 3A_i(x, x', x''),$$

$$\overset{1}{E}_i = G_{ij} x''^{jj} + \Gamma_i(x, x', x''),$$

where

$$G_{ij} = 3A_{i(1)j} + A_{j(1)i} - 2B_{(2)i(2)j},$$

$$A_i = A_{i(0)j} x'^j - \frac{1}{3}B_{(2)i},$$

$$\Gamma_i = B_{(1)i} - 2B_{(2)i(1)j} x''^{jj} - 2B_{(2)i(0)j} x'^j + 3(A_{i(1)j(1)k} x''^{jj} x''^{kk} + 2A_{i(1)j(0)k} x''^{jj} x''^{kk} + A_{i(0)j(0)k} x''^{kk} x'^j + A_{i(0)j} x''^{jj}).$$

Then we have

$$\frac{1}{3}A^{ik} \overset{2}{E}_i = x''^{kk} + A^k, \quad A^k = A^{ik} A_i,$$

$$G^{ik} \overset{1}{E}_i = x''^{kk} + \Gamma^k, \quad \Gamma^k = G^{ik} \Gamma_i,$$

$$A^{ik} A_{i(1)j} = \delta_j^k, \quad G^{ik} G_{ik} = \delta_j^k.$$

By use of A^i and Γ^i , we shall introduce the linear connection which is invariant for a change of parameter t .

§1. Synge vector

In the space we have *Zermelo's conditions*

$$\sum_{\beta=p}^3 \beta x^{(\beta-p+1)t} F_{(\beta)t} = p \delta_p^j F, \quad (\rho=1, 2, 3),$$

that is,

$$(1.1) \quad A_i x'^i = 0, \quad A_{i(1)j} x'^j = (p-3)A_i,$$

$$B_{(2)i} x'^i + 3A_i x''^{ii} = 0, \quad B_{(1)i} x'^i + 2B_{(2)i} x''^{ii} = pB.$$

From (1.1) we have

$$(1.2) \quad A_{i(1)j} x'^i = -A^j, \quad A_{i(0)j} x'^i = 0, \quad B_{(2)i(2)j} x'^i = -3A_j.$$

Furthermore, using (1.1) and (1.2), after some calculation we have

$$(1.3) \quad \Gamma_i x'' = pB,$$

$$(1.4) \quad \Gamma_{i(2)j} x'^j = -3G_{ik} x''^k + (2p-3)\overset{2}{E}_i,$$

$$(1.5) \quad \Gamma_{i(2)j} x'^j = -3x''^i + (2p-3)\overset{2}{E}^i, \quad \overset{2}{E}^i = G^j \overset{2}{E}_j,$$

$$(1.6) \quad \overset{2}{E}_{(2)j} x'^j = \frac{3(p-2)}{3p-4} x''^i,$$

$$(1.7) \quad A_{(2)j}^i x'^j = \frac{1}{p-3} x''^i.$$

Now, for a change of parameter t , we have

$$(1.8) \quad A_i(\bar{t}) = \gamma^{p-3} A_i, \quad A_{i(1)j}(\bar{t}) = \gamma^{p-4} A_{i(1)j}, \quad A_{i(1)j(1)k}(\bar{t}) = \gamma^{p-5} A_{i(1)j(1)k}, \\ B_{(2)l(2)j}(\bar{t}) = \gamma^{p-4} B_{(2)l(2)j}, \quad G_{ij}(\bar{t}) = \gamma^{p-4} G_{ij}, \quad G^{ij}(\bar{t}) = \gamma^{-p+4} G^{ij}.$$

Next, using (1.8) we obtain

$$(1.9) \quad \overset{2}{E}_i = \gamma^{p-2} \overset{2}{E}_i + \gamma^{p-4} \frac{d\gamma}{dt} 3(p-2) A_i,$$

$$(1.10) \quad \overset{1}{E}_i = \gamma^{p-1} \overset{1}{E}_i + S \gamma^{p-3} \frac{d\gamma}{dt} \overset{2}{E}_i - \left\{ \gamma^{p-4} \frac{d^2\gamma}{dt^2} Q - \gamma^{p-5} \left(\frac{d\gamma}{dt} \right)^2 R \right\} \overset{3}{E}_i$$

where

$$S = 2p-3, \quad Q = -3p+4, \quad R = 3(p^2-5p+5).$$

Then, using (1.8) and (1.9) we have

$$(1.11) \quad K(\bar{t}) = \gamma K + \gamma^{-1} \frac{d\gamma}{dt}$$

where

$$K = G^{ij} G^{jk} A_{i(1)j(1)k} \overset{2}{E}_i / \frac{3(p-2)(p-4)}{3p-4} G^{ij} A_{i(1)j},$$

or

$$K = G^{ij} A^{jk} A_{i(1)j(1)k} \overset{2}{E}_i / \frac{3(p-2)(p-4)}{p-3} G^{ij} A_{i(1)j}.$$

At last, from (1.10), we obtain

$$(1.12) \quad \Gamma^j(\bar{t}) = \gamma^3 \Gamma^j + \gamma^3 x^{(3)j} - x^{(3)j} + \gamma \frac{d\gamma}{dt} S G^{ij} \overset{2}{E}_i - \left\{ \gamma^{p-4} \frac{d^2\gamma}{dt^2} Q - \gamma^{p-5} \left(\frac{d\gamma}{dt} \right)^2 R \right\} G^{ij} \overset{3}{E}_i,$$

$$(1.13) \quad \overset{2}{E}^i = \gamma^2 \overset{2}{E}^i + \frac{3(p-2)}{3p-4} x''^i \frac{d\gamma}{dt}.$$

Differentiating (1.12) and (1.13) with respect to $x^{(2)k}$, we have

$$(1.14) \quad \Gamma_{(2)k}^i(\bar{t}) = \gamma \Gamma_{(2)k}^i + \Gamma_{(2)j(2)k}^i x''^j \gamma^{-1} \frac{d\gamma}{dt},$$

$$(1.15) \quad \overset{2}{E}_{(2)k}^i(\bar{t}) = \overset{2}{E}_{(2)k}^i.$$

§ 2. Connection parameter (I)

Since $1/3 A^{ik} \overset{2}{E}_i$ is the vector, we have

$$\bar{x}'''^i + \bar{\Gamma}^i(x, x', x'') = \frac{\partial \bar{x}^i}{\partial x^k} (x'''^k + \Gamma^k), \quad \bar{x}''''^i + \bar{\Gamma}^i(x, x', x'') = \frac{\partial \bar{x}^i}{\partial x^k} (x''''^k + \Gamma^k)$$

for the coordinate transformation.

Hence we have

$$(2.1) \quad \bar{A}^t = \frac{\partial \bar{x}^t}{\partial x^k} A^k + \frac{\partial \bar{x}^t}{\partial x^k} x''^k - x'''^t,$$

$$(2.2) \quad \bar{\Gamma}^t = \frac{\partial \bar{x}^t}{\partial x^k} \Gamma^k + \frac{\partial \bar{x}^t}{\partial x^k} x'''^k - x''''^t.$$

Differentiating (2.1) and (2.2) with respect to x'^j and x''^j respectively, we have

$$(2.3) \quad \bar{\Gamma}^t_{(2)j} = \frac{\partial \bar{x}^t}{\partial x^k} \frac{\partial x^t}{\partial x'^j} \Gamma^k_{(2)t} + 3 \frac{\partial \bar{x}^t}{\partial x^k} \frac{\partial^2 x^k}{\partial x^j \partial x^n} \bar{x}''^n,$$

$$(2.4) \quad \bar{A}^t_{(1)j} = \frac{\partial \bar{x}^t}{\partial x^k} \frac{\partial x^t}{\partial x'^j} A^k_{(1)t} + 2 \frac{\partial \bar{x}^t}{\partial x^k} \left(\frac{\partial x^k}{\partial x'^j} \right)^{(1)} + \frac{\partial \bar{x}^t}{\partial x^k} \frac{\partial x''^k}{\partial x'^j} A^k_{(2)t}.$$

Substituting (2.3) to (2.4) we have

$$\left(\bar{A}^t_{(1)j} - \frac{2}{3} \bar{\Gamma}^t_{(2)j} \right) \frac{\partial x^k}{\partial x^t} = \frac{\partial x^t}{\partial x'^j} M^t_k + \frac{\partial \bar{x}^t}{\partial x^t} \frac{\partial x''^p}{\partial x'^j} A^t_{(2)p},$$

where

$$M^t_j = M^t_{(1)j} - \frac{2}{3} \Gamma^t_{(2)j}.$$

From $A^t_{(2)p} \cong 0$, we have by (1.7)

$$(2.5) \quad \Omega^m_k \Omega_l^k = \delta_l^m, \quad \Omega^k_l \Omega_k^j = \delta_l^j, \quad \Omega^m_k A_m = -A_k, \quad \Omega^m_k = A^m_{(2)k},$$

$$(2.6) \quad \frac{\partial x^m}{\partial x^t} \bar{M}^t_j \Omega_m^t = \frac{\partial x^k}{\partial x'^j} M^m_k \Omega_m^t + 2 \frac{\partial^2 x^t}{\partial x^j \partial x^n} \bar{x}''^n.$$

On the other hand, differentiating (2.3) with respect to x'^k we have

$$\begin{aligned} \bar{\Gamma}^t_{(2)j(1)k} &= \frac{\partial \bar{x}^t}{\partial x^m} \frac{\partial x^t}{\partial x'^j} \Gamma^m_{(2)t(1)q} \frac{\partial x^q}{\partial x'^k} + 3 \frac{\partial \bar{x}^t}{\partial x^p} \frac{\partial^2 x^p}{\partial x^j \partial x^k} + \frac{\partial \bar{x}^t}{\partial x^m} \frac{\partial x^t}{\partial x'^j} \frac{\partial x''^q}{\partial x'^k} \Gamma^m_{(2)t(2)q} \\ &= \frac{\partial \bar{x}^t}{\partial x^m} \frac{\partial x^t}{\partial x'^j} \frac{\partial x^q}{\partial x'^k} \Gamma^m_{(2)t(1)q} + 3 \frac{\partial \bar{x}^t}{\partial x^p} \frac{\partial^2 x^p}{\partial x^j \partial x^k} \\ &\quad + \frac{\partial \bar{x}^t}{\partial x^m} \frac{\partial x^t}{\partial x'^j} \Gamma^m_{(2)t(2)q} \left(\bar{M}^t_k \Omega_r^q \frac{\partial x^r}{\partial x^t} - M^t_r \Omega_t^q \frac{\partial x^r}{\partial x^k} \right). \end{aligned}$$

Seeing that $\Gamma^t_{(2)j(2)k}$ is a tensor, we have

$$(2.7) \quad \bar{\Pi}^t_{jk} = \frac{\partial x^p}{\partial x'^j} \frac{\partial x^q}{\partial x'^k} \frac{\partial \bar{x}^t}{\partial x^t} \Pi^t_{pq} - \frac{\partial x^p}{\partial x'^j} \frac{\partial x^q}{\partial x'^k} \frac{\partial^2 \bar{x}^t}{\partial x^p \partial x^q},$$

where

$$3 \Pi^t_{jk}(x, x', x'') = \Gamma^t_{(2)j(1)k} - \Omega^q_p M^p_k \Gamma^t_{(2)j(2)q}.$$

Thus, using (2.7) we can define the covariant differential

$$\delta v^t = dv^t + \Pi^t_{jk} v^j dx^k.$$

The base connection is given by

$$\begin{aligned} \delta x''^t &= dx''^t + \Pi^t_{jk} x'^j dx^k, & x^t &= x'''^t + \Pi^t_{jk} x'^j x'^k, \\ \delta x^t &= dx^t + \Pi^t_{jk} x^j dx^k, & x^t &= x'^t + \Pi^t_{jk} x^j x'^k. \end{aligned}$$

§ 3. Connection parameter (II)

Substituting (2.3) to the third term in the right member of (2.4), we have

$$(3.1) \quad \bar{A}^t_j = \frac{\partial x^p}{\partial x'^j} \frac{\partial \bar{x}^t}{\partial x^q} A^q_p - 2 \frac{\partial x^t}{\partial x'^j} \frac{\partial^2 \bar{x}^t}{\partial x^t \partial x^p} x''^p,$$

or

$$(3.1)' \quad \frac{\partial \bar{x}^t}{\partial x^k} A_j^t = \frac{\partial \bar{x}^t}{\partial x^q} A_k^q - 2 \frac{\partial^2 \bar{x}^t}{\partial x^k \partial x^p} x'^p,$$

where

$$(3.2) \quad A_j^t = A_{(1)j}^t - \frac{2}{3} \Gamma_{(2)j}^p A_{(2)p}^t.$$

After short calculation we obtain

$$(3.3) \quad A_j^t x'^j = 2A^t - \frac{2(2p-3)}{3} E^p A_{(2)p}^t = 2A^t - \frac{2(2p-3)}{3} A_{(2)j}^t x''^j.$$

Differentiating (3.1) with respect to x'^k and using (3.1), we have

$$(3.4) \quad \bar{A}_{jk}^t = \frac{\partial x^p}{\partial \bar{x}^j} \frac{\partial \bar{x}^t}{\partial x^q} \frac{\partial x^t}{\partial \bar{x}^k} A_{pt}^q - 2 \frac{\partial x^t}{\partial \bar{x}^k} \frac{\partial x^m}{\partial \bar{x}^j} \frac{\partial^2 \bar{x}^t}{\partial x^m \partial x^l},$$

or

$$(3.4)' \quad \bar{A}_{jk}^t \frac{\partial \bar{x}^j}{\partial x^t} \frac{\partial \bar{x}^k}{\partial x^s} = \frac{\partial \bar{x}^t}{\partial x^q} A_{ts}^q - 2 \frac{\partial^2 \bar{x}^t}{\partial x^t \partial x^s},$$

where

$$(3.5) \quad A_{jk}^t = A_{(1)jk}^t - A_{j(2)s}^t A_k^s.$$

Thus, using (3.4) we can define the covariant differential, but we will modify.

Now, if we put

$$*A_j^t = A_j^t + \frac{2(2p-3)}{q(p-4)} A_{(1)j}^{pt} B_{(2)p(2)q} E^q,$$

then the transformation of $*A_j^t$, for the coordinate transformation, is (3.1), because the second term in the above equation is a tensor. Using (3.3) we have $*A_j^t x'^j = 2A^t$. Hence, by means of the process which obtained (3.4), $*A_{jk}^t = *A_{j(1)k}^t - *A_{j(2)s}^t *A_k^s$ have the same transformation law as the case of A_{jk}^t .

Thus we can define the covariant differential

$$\delta v^t = dv^t + *A_{jk}^t v^j dx^k.$$

§ 4. Connection parameter (III)

In § 2 and § 3, by use of A^k and Γ^k we have introduced the connection parameter Π_{jk}^t , $*A_{jk}^t$. In this paragraph, from only Γ^t we shall introduce a connection parameter which is not intrinsic. Differentiating (2.3) with respect to x'^k we have

$$\bar{\Gamma}_{(2)j(1)k}^t = \frac{\partial x^p}{\partial \bar{x}^j} \frac{\partial x^q}{\partial \bar{x}^k} \left(\frac{\partial \bar{x}^t}{\partial x^l} \Gamma_{(2)p(1)q}^l - \frac{\partial^2 \bar{x}^t}{\partial x^p \partial x^q} \right) + \frac{\partial x^p}{\partial \bar{x}^j} \frac{\partial \bar{x}^t}{\partial x^l} \frac{\partial x'^q}{\partial \bar{x}^k} \Gamma_{(2)p(2)q}^l.$$

Since $\Gamma_{(2)p(2)q}^t$ is a tensor, using (2.3) we have

$$(4.1) \quad \Gamma_{jk}^t = \frac{\partial x^p}{\partial \bar{x}^j} \frac{\partial x^q}{\partial \bar{x}^k} \frac{\partial \bar{x}^t}{\partial x^l} \Gamma_{pq}^l - \frac{\partial x^p}{\partial \bar{x}^q} \frac{\partial x^q}{\partial \bar{x}^k} \frac{\partial^2 \bar{x}^t}{\partial x^p \partial x^q},$$

where

$$(4.2) \quad 3\Gamma_{jk}^t = \Gamma_{(2)j(1)k}^t - \frac{2}{3} \Gamma_{(2)j(2)l}^t \Gamma_{(2)l}^t.$$

Thus, from (4.1), we define The covariant differentiation

$$(4.3) \quad \delta v^i = dv^i + \Gamma_{jk}^i v^j dx^k$$

and the base connection

$$(4.4) \quad \begin{aligned} \delta x^i &= dx^i + \Gamma_{jk}^i x'^j dx^k, & x^i &= x'^i + \Gamma_{jk}^i x'^j x'^k, \\ \delta x^i &= dx^i + \Gamma_{jk}^{(2)i} x'^j dx^k, & x^i &= x'^i + \Gamma_{jk}^{(2)i} x'^j x'^k. \end{aligned}$$

§ 5. Intrinsic connection parameter

In this paragraph we shall introduce the intrinsic connection parameter.

Theorem 5.1. If we define

$$* \prod_{jk}^i = \Gamma_{jk}^i - \Omega_v^u M_k^v \Gamma_{j(2)u}^i, \quad \Gamma_{jk}^i = \Gamma_{j(1)k}^i + K \Gamma_{j(2)k}^i,$$

$$\Gamma_j^i = \frac{1}{3} (\Gamma_{(2)j}^i - K \Gamma_{(2)k(2)j}^i x'^k), \quad N_j^i = A_j^i - 2\Gamma_j^i,$$

$$A_j^i = A_{(1)j}^i - \frac{1}{p-3} (\delta_j^i - A_{(1)j}^i) K,$$

then $* \prod_{jk}^i$ is the intrinsic connection parameter.

Proof: From (1.11) and (1.14), we have

$$(5.1) \quad \Gamma_j^i(\bar{t}) = r \Gamma_j^i$$

for a change of parameter and

$$(5.2) \quad \bar{\Gamma}_j^i = \frac{\partial x^p}{\partial x^j} \frac{\partial \bar{x}^i}{\partial x^q} \Gamma_{pq}^i - \frac{\partial x^q}{\partial x^j} \frac{\partial^2 \bar{x}^i}{\partial x^p \partial x^q} x'^p$$

for the coordinate transformation. Differentiating (5.1) with respect to x'^k and using (1.11), we have

$$(5.3) \quad \Gamma_{j(1)k}^i(\bar{t}) + K(\bar{t}) \Gamma_{j(2)k}^i(\bar{t}) = \Gamma_{j(1)k}^i + K \Gamma_{j(2)k}^i$$

i.e., $\Gamma_{jk}^i(\bar{t}) = \Gamma_{jk}^i$.

Further, differentiating (5.2) with respect to x'^k we have

$$\bar{\Gamma}_{j(1)k}^i = \frac{\partial x^q}{\partial x^k} \frac{\partial x^p}{\partial x^j} \frac{\partial \bar{x}^i}{\partial x^t} \Gamma_{p(1)q}^i - \frac{\partial x^q}{\partial x^k} \frac{\partial x^t}{\partial x^j} \frac{\partial^2 \bar{x}^i}{\partial x^t \partial x^q} + 2 \frac{\partial x^p}{\partial x^j} \frac{\partial \bar{x}^i}{\partial x^t} \Gamma_{p(2)q}^i \frac{\partial^2 x^q}{\partial x^k \partial x^t} \bar{x}'^t.$$

Since $K \Gamma_{j(2)k}^i$ is a tensor, we obtain

$$(5.4) \quad \bar{\Gamma}_{jk}^i = \frac{\partial x^q}{\partial x^k} \frac{\partial x^p}{\partial x^j} \frac{\partial \bar{x}^i}{\partial x^t} \Gamma_{pq}^i - \frac{\partial x^q}{\partial x^k} \frac{\partial x^t}{\partial x^j} \frac{\partial^2 \bar{x}^i}{\partial x^t \partial x^q} + 2 \frac{\partial x^p}{\partial x^j} \frac{\partial \bar{x}^i}{\partial x^t} \Gamma_{p(2)q}^i \frac{\partial^2 x^q}{\partial x^k \partial x^t} \bar{x}'^t.$$

On the other hand, we have

$$\bar{A}_j^i = \frac{\partial x^q}{\partial x^j} \frac{\partial \bar{x}^i}{\partial x^p} A_p^q - 2 \frac{\partial x^k}{\partial x^j} \frac{\partial^2 \bar{x}^i}{\partial x^k \partial x^t} x'^t + 2 \frac{\partial^2 x^t}{\partial x^j \partial x^n} \frac{\partial \bar{x}^i}{\partial x^q} \bar{x}'^n A_{(2)t}^i.$$

Substituting (5.2) to the above equation we have

$$(5.5) \quad \bar{N}_j^i = \frac{\partial x^p}{\partial x^j} \frac{\partial \bar{x}^i}{\partial x^q} N_p^q + 2 \frac{\partial^2 x^t}{\partial x^j \partial x^n} \bar{x}'^n \frac{\partial \bar{x}^i}{\partial x^q} A_{(2)t}^i.$$

Multiplying (5.5) with Ω_q^t and summing with q we have

$$(5.6) \quad \bar{N}_j^p \Omega_q^t \frac{\partial x^q}{\partial x^p} = \frac{\partial x^p}{\partial x^j} N_p^q \Omega_q^t + 2 \frac{\partial^2 x^t}{\partial x^j \partial x^n} \bar{x}'^n.$$

Thus, substituting (5.6) to (5.4) we obtain

$$(5.7) \quad {}^* \bar{\Pi}_1^i{}_{jk} = \frac{\partial x^p}{\partial \bar{x}^j} \frac{\partial x^q}{\partial \bar{x}^k} \frac{\partial \bar{x}^t}{\partial x^t} {}^* \Pi_1^{t pq} - \frac{\partial x^p}{\partial \bar{x}^j} \frac{\partial x^q}{\partial \bar{x}^k} \frac{\partial^2 \bar{x}^i}{\partial x^p \partial x^q}$$

Further, from

$$\Omega_v^u(\bar{t}) = \Omega_v^u, \quad N_j^i(\bar{t}) = \gamma N_j^i, \quad \Gamma_{j(2)k}^i(\bar{t}) = \gamma^{-1} \Gamma_{j(2)k}^i$$

we have

$${}^* \bar{\Pi}_1^i{}_{jk}(\bar{t}) = {}^* \Pi_1^i{}_{jk}(t) \quad \text{Q.E.D.}$$

By means of Theorem 5.1 we can define the intrinsic covariant differential

$$(5.8) \quad \delta v^i = dv^i + {}^* \Pi_1^i{}_{jk} v^j dx^k.$$

Theorem 5.2: If we define

$$2 {}^* \Pi_2^i{}_{jk} = {}^* A_{1jk}^i - {}^* A_{1j(2)s}^i {}^* A_{1k}^s, \quad {}^* A_{1jk}^i = {}^* A_{j(1)k}^i + K {}^* A_{1j(2)k}^i, \quad {}^* A_{1j}^i = A_{1j}^i - 2\Gamma_{1j}^p A_{(2)p}^i$$

then ${}^* \Pi_2^i{}_{jk}$ is an intrinsic connection parameter.

Proof: Both A_j^i and $A_{(1)j}^i$ are given the same transformation, and hence we have

$$(5.9) \quad \bar{A}_1^i = \frac{\partial x^q}{\partial \bar{x}^j} \frac{\partial \bar{x}^j}{\partial x^p} A_{1q}^p - 2 \frac{\partial x^k}{\partial \bar{x}^j} \frac{\partial^2 \bar{x}^i}{\partial x^k \partial x^t} x'^t + 2 \frac{\partial^2 \bar{x}^i}{\partial \bar{x}^j \partial \bar{x}^s} \bar{x}'^s \frac{\partial \bar{x}^t}{\partial x^q} A_{(2)q}^t$$

for the coordinate transformation and

$$(5.10) \quad \bar{A}_1^i(\bar{t}) = \gamma A_1^i(t)$$

for a change of parameter. From (5.2) and (5.9), we have

$$(5.11) \quad {}^* \bar{A}_1^i \frac{\partial \bar{x}^j}{\partial x^k} = \frac{\partial x^p}{\partial \bar{x}^j} \frac{\partial \bar{x}^t}{\partial x^q} {}^* A_{1p}^q - 2 \frac{\partial^2 \bar{x}^i}{\partial x^p \partial x^t} \bar{x}'^t$$

From (5.10) we obtain

$$(5.12) \quad {}^* A_{1j(2)k}^i = \gamma^{-1} {}^* A_{1j(2)k}^i, \quad {}^* A_{j(1)k}^i + K {}^* A_{1j(2)k}^i = {}^* A_{j(1)k}^i + K {}^* A_{1j(2)k}^i$$

and, since $K {}^* A_{1j(2)k}^i$ is a tensor, we have

$${}^* \bar{A}_1^i = \frac{\partial x^p}{\partial \bar{x}^j} \frac{\partial \bar{x}^t}{\partial x^q} \frac{\partial x^t}{\partial \bar{x}^k} {}^* A_{1pt}^q - 2 \frac{\partial x^m}{\partial \bar{x}^k} \frac{\partial x^t}{\partial \bar{x}^j} \frac{\partial^2 \bar{x}^i}{\partial x^t \partial x^m} + 2 \frac{\partial x^p}{\partial \bar{x}^j} \frac{\partial x^t}{\partial \bar{x}^q} \frac{\partial^2 \bar{x}^i}{\partial x^s \partial \bar{x}^k} \bar{x}'^s {}^* A_{p(2)t}^q$$

At last, substituting (5.11) to the above equations, we have

$${}^* \bar{\Pi}_2^i{}_{jk} = \frac{\partial x^p}{\partial \bar{x}^j} \frac{\partial x^t}{\partial \bar{x}^k} \frac{\partial \bar{x}^t}{\partial x^q} {}^* \Pi_2^{i pq} - 2 \frac{\partial x^p}{\partial \bar{x}^k} \frac{\partial x^q}{\partial \bar{x}^j} \frac{\partial^2 \bar{x}^i}{\partial x^p \partial x^q}$$

Further, from (5.12) and (5.10), we have

$${}^* \bar{\Pi}_2^i{}_{jk}(\bar{t}) = {}^* \Pi_2^i{}_{jk}(t) \quad \text{Q.E.D.}$$

Hence we obtain the covariant differential

$$\delta v^i = dv^i + {}^* \Pi_2^i{}_{jk} v^j dx^k.$$

Theorem 5.3 If we define

$${}^* \Pi_3^i{}_{jk} = \Gamma_{j(1)k}^i + K \Gamma_{j(2)k}^i - \frac{2}{3} \Gamma_{j(2)t}^i \Gamma_k^t$$

then ${}^* \Pi_3^i{}_{jk}$ is an intrinsic connection parameter.

Proof: Differentiating (5.2) with respect to x'^k and using (5.2), we have

$$(5.13) \quad \bar{\Gamma}_{j(1)k}^i - 2\bar{\Gamma}_k^p \bar{\Gamma}_{j(2)p}^i = \frac{\partial x^q}{\partial \bar{x}^k} \frac{\partial x^p}{\partial \bar{x}^j} \frac{\partial \bar{x}^t}{\partial x^t} \left(\Gamma_{p(1)q}^t - 2\Gamma_q^s \Gamma_{p(2)s}^t \right) - \frac{\partial x^q}{\partial \bar{x}^k} \frac{\partial x^t}{\partial \bar{x}^j} \frac{\partial^2 \bar{x}^i}{\partial x^t \partial x^q}$$

Since $\Gamma_k^p \Gamma_{j(2)p}^i$ is an intrinsic tensor, making use of (1.11) we have

$$\begin{aligned} (\Gamma_{j(1)k}^i - 2\Gamma_k^p \Gamma_{j(2)p}^i)(\bar{t}) &= \Gamma_{j(1)k}^i - 2\Gamma_k^p \Gamma_{j(2)p}^i - \gamma^{-2} \frac{d\gamma}{dt} \Gamma_{j(2)k}^i \\ &= \Gamma_{j(1)k}^i - 2\Gamma_k^p \Gamma_{j(2)p}^i - (K(\bar{t}) - \gamma K) \gamma^{-1} \Gamma_{j(2)k}^i, \end{aligned}$$

and consequently

$$(5.14) \quad * \Pi_{jk}^i(\bar{t}) = * \Pi_{jk}^i(t)$$

Thus, since $K \Gamma_{j(2)k}^i$ is a tensor, by use of (5.13) and (5.14), we have

$$* \bar{\Pi}_{jk}^i = \frac{\partial x^p}{\partial x^j} \frac{\partial x^q}{\partial x^k} \frac{\partial \bar{x}^i}{\partial x^l} * \Pi_{pq}^l - \frac{\partial x^p}{\partial x^j} \frac{\partial x^q}{\partial x^k} \frac{\partial^2 \bar{x}^i}{\partial x^p \partial x^q} \quad \text{Q.E.D.}$$

§ 6. Base connection

We shall introduce *the base connection* in the special Kawaguchi space.

The expressions

$$(6.1) \quad \delta_{\alpha}^{(1)} x^i = d(x^i/F)^* + \Pi_{jk}^i(x^j/F) dx^k, \quad \delta_{\alpha}^{(2)} x^i = dx^i + * \Pi_{jk}^i(x^j) dx^k, \quad \alpha = 1, 2, 3,$$

are the components of a contravariant vector. The expressions (6.1) give the base connection in K_n^2 , a special Kawaguchi space of order 2 and dimension n.

Along the base curve the equations (6.1) become

$$(6.2) \quad \delta_{\alpha}^{(2)} x^i = \frac{d}{dt} (x^i/F) + * \Pi_{jk}^i x^j x^k F^{-1}, \quad x^i = \frac{d}{dt} x^i + * \Pi_{jk}^i x^j x^k.$$

Thus, making use of (6.1) and (5.2), we have the covariant derivative.

Lately the author have introduced the intrinsic covariant differential in the special Kawaguchi space with arc length

$$s = \int \{A_i(x, x', \dots, x^{(m-2)t}) x^{(m)t} + B(x, x', \dots, x^{(m-1)t})\}^{1/p} dt \quad [4].$$

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