



計量 $s = \int (A_{ix}''^I + B)^{<1/p>} dt$ をもつ n 次元空間における曲線の共形論

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On the Conformal Theory of Curves in an n -Dimensional Space with the Arc Length $s = \int (A_i x''^i + B)^{1/p} dt$

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叶 長太郎：計量 $s = \int (A_i x''^i + B)^{1/p} dt$ をもつ
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Introduction. The fundamental theory of the n -dimensional Kawaguchi space with the metrics $s = \int F^{1/p} dt$, where $F = A_i x''^i + B$, was developed by Prof. A. Kawaguchi [1], [2]. The general curve theory of the space was studied by S. Ide [3], [4]. In the previous paper [5] the present author has introduced a conformal theory in such a space.

In the present paper we shall develop a conformal curve theory in the space.

§ I. Conformal parameter. In the previous paper we have seen the following relations between the fundamental quantities:

$$(1.1) \quad \begin{aligned} \bar{A}_i \left(x, \frac{dx}{dt} \right) &= e^{\nu\sigma} A_i \left(x, \frac{dx}{dt} \right), \\ \bar{A}_{i(j)} &= e^{\nu\sigma} (A_{i(j)} + p\sigma_{(j)} A_i), \\ \bar{G}_{i,j} &= e^{\nu\sigma} \{ G_{i,j} + p(2\sigma_{(j)} A_i - \sigma_{(i)} A_j) \}, \\ \bar{A}_{ij} &= e^{\nu\sigma} (A_{ij} + p\sigma_j A_i), \end{aligned}$$

where the parameter t is a conformal parameter defined along the curve.

Instead of the parameter t , if we take the invariant parameter s defined as the arc length of a curve, then (1.1) becomes

$$(1.2) \quad \begin{aligned} \bar{A}_i \left(x, \frac{dx}{ds} \right) &= e^{2\sigma} A_i \left(x, \frac{dx}{ds} \right), \quad \bar{B} \left(x, \frac{dx}{ds} \right) = B \left(x, \frac{dx}{ds} \right), \\ \bar{A}_{i(j)} &= e^{3\sigma} (A_{i(j)} + p\sigma_{(j)} A_i), \\ \bar{A}_{ij} &= e^{2\sigma} (A_{ij} + p\sigma_j A_i), \\ \bar{G}_{i,j} &= e^{3\sigma} (G_{i,j} + p(2\sigma_{(j)} A_i - \sigma_{(i)} A_j)), \end{aligned}$$

Then the conformal base connection

$$x^{[2]i} = \frac{d^2 x^i}{ds^2} + 2H^i \left(x, \frac{dx}{ds} \right),$$

where $H^i = \Gamma^i - Kx'^i$ and $K = G^{ij} \nabla_j A_i \left(\frac{p^2 - 3}{p(2p - 3)} - G^{ij} A_{i(j)} \right)$, are not conformal. Hence we have from (1.2) the following transformation law under the conformal transformation

$$(1.3) \quad \bar{x}^{[2]i} = e^{-2\sigma} \left(x^{[2]i} - \frac{d\sigma}{ds} \cdot \frac{dx^i}{ds} \right).$$

Then we find, for the conformal differential of $x^{[2]i}$ along the curve $x^i = x^i(s)$,

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$$\frac{\partial^* \tilde{x}^{[2]i}}{d\tilde{s}} = e^{-3\sigma} \left\{ \frac{\partial^* x^{[2]i}}{ds} + \left(2 \left(\frac{d\sigma}{ds} \right)^2 - \frac{d^2\sigma}{ds^2} \right) \frac{dx^i}{ds} - 3 \frac{d\sigma}{ds} x^{[2]i} \right\}.$$

If these equations be multiplied by A_i and summed for i , we obtain

$$(1.4) \quad \tilde{Q}(\tilde{s}) = e^{-\sigma} \left\{ Q(s) - \frac{d\sigma}{ds} \right\},$$

where

$$Q(s) = \frac{1}{2} A_i \frac{\partial^* x^{[2]i}}{ds} = -\frac{1}{3} \{ G_{ij} x^{[2]i} x^{[2]j} + (2p-3)K \}.$$

Thus if we put

$$\varphi(t) = \left[\{t, s\} - \frac{1}{2} \left(2 \frac{dQ}{ds} - Q^2 \right) \right] / \left(\frac{dt}{ds} \right)^2,$$

where $\{t, s\}$ denotes the Schwarzian derivative of t with respect to s , then from (1.4) it follows that the quantity $\varphi(t)$ is invariant under the conformal transformation, and consequently the parameter t defined by the equation $\varphi(t)=0$ is of conformal character.

§ 2. Conformal arc length. If we put $x^{*[2]i} = x^{[2]i} - Q(s) \frac{dx^i}{ds}$, then by means of (1.3) and (1.4) the transformation law of $x^{*[2]i}$ under the conformal transformation is $\tilde{x}^{*[2]i} = e^{-2\sigma} x^{*[2]i}$. Further if we put $V^i = \frac{\partial^* x^{*[2]i}}{ds} - 2Q(s) \frac{dx^i}{ds}$, we have $\tilde{V}^i = e^{-3\sigma} V^i$, where

$$(2.1) \quad V^i = \frac{\partial x^{[2]i}}{ds} + \left\{ -K_{(j)} x^{[2]j} - \frac{1}{3} \frac{\partial \nu}{ds} + \frac{2}{9} R\nu \right\} \frac{dx^i}{ds} - R x^{[2]i}$$

$$\frac{\partial x^{[2]i}}{ds} = \frac{dx^{[2]i}}{ds} + \Gamma_{(j)}^i x^{[2]j}, \quad R = A_i x^{[3]i},$$

$$\nu = R + 3K.$$

Then we have $A_i V^i = 0$.

In parallel with the above-mentioned equations, covariant representation of them is that

$$(2.2) \quad \tilde{S}_i^{[1]} = e^\sigma S_i^{[1]}, \quad \tilde{S}_i^{[2]} = S_i^{[2]}$$

$$S_i^{[1]} = \frac{\partial^* A_i}{ds} + \frac{2}{3} (R - 3K) A_i = A_i^{[1]} + \left(\frac{2}{3} R - K \right) A_i,$$

$$S_i^{[2]} = \frac{\partial^* A_i^{[1]}}{ds} + \frac{1}{3} (R - 3K) S_i^{[1]} = A_i^{[1]} + (R - K) A_i^{[1]} + T A_i - K_{(i)},$$

$$T = \frac{d}{ds} \left(\frac{2}{3} R - K \right) + \frac{1}{3} R \left(\frac{2}{3} R - K \right),$$

$$S_i^{[1]} x'^i = -1, \quad S_i^{[2]} x'^i = 0.$$

Now we shall assume that $V^i \neq 0$.

Thus from (2.1) and (2.2) the quantity $J(s)$ defined by

$$(2.3) \quad J(s) = (S_i^{[2]} V^i)^{1/3}$$

$$= (-k_2 \bar{k}_2 - k_1 K + R K_{(i)} x^{[2]i} - K_{(i)} x^{[3]i})^{1/3}, \quad [3]$$

is transformed into $\tilde{J}(\tilde{s}) = e^{-\sigma} J(s)$.

Hence we can see that the parameter $\kappa(s)$ defined by

$$(2.4) \quad \kappa(s) = \int J(s) ds$$

is a conformal parameter defined along the curve. We shall call it conformal arc length of the curve.

§ 3. Conformal covariant derivative and conformal Frenet's Formulas.

Let us consider a curve $x^i = x^i(\kappa)$ and assume that κ is a conformal parameter. If v^i is a relative conformal contravariant vector defined along the curve, that is, $\tilde{v}^i = e^{-\sigma} v^i$, then from (1.4) we obtain

$$\frac{Dv^i}{d\kappa} = \frac{\partial^{**} v^i}{d\kappa} + J^{-1} Q v^i$$

that the transformation law of $Dv^i/d\kappa$ under conformal transformation is $D\tilde{v}^i/d\kappa = e^{-\sigma} Dv^i/d\kappa$. We shall call $Dv^i/d\kappa$ the conformal covariant derivative of the relative conformal contravariant vector v^i along the curve.

For a relative conformal covariant vector we can define conformal covariant derivative in the same way, that is,

$$\frac{Dv_i}{d\kappa} = \frac{\partial^{**} v_i}{d\kappa} - J^{-1} Q v_i$$

Now if we put $\overset{1}{\gamma}^i = \frac{dx^i}{ds}$ and $\lambda_i = J A_i$, then $\overset{1}{\gamma}^i$ and λ_i is transformed into $\overset{1}{\tilde{\gamma}}^i = e^{-\sigma} \overset{1}{\gamma}^i$, and $\overset{1}{\tilde{\lambda}}_i = e^{\sigma} \lambda_i$ respectively. Then we have $\lambda_i D\overset{1}{\gamma}^i/d\kappa = 1$, $\frac{D\lambda_i}{d\kappa} \overset{1}{\gamma}^i = -1$ and $\lambda_i \overset{1}{\tilde{\gamma}}^i = 0$.

Thus with the aid of the following relations:

$$(3.1) \quad \begin{aligned} \lambda_i \overset{1}{\tilde{\gamma}}^i &= 0, & \lambda_i \overset{2}{\tilde{\gamma}}^i &= 1, & \lambda_i \overset{3}{\tilde{\gamma}}^i &= 0, & \dots & \dots, & \lambda_i \overset{n}{\tilde{\gamma}}^i &= 0, \\ \lambda_i \overset{1}{\tilde{\gamma}}^i &= -1, & \lambda_i \overset{2}{\tilde{\gamma}}^i &= k^*, & \lambda_i \overset{3}{\tilde{\gamma}}^i &= 0, & \dots & \dots, & \lambda_i \overset{n}{\tilde{\gamma}}^i &= 0, \\ & \vdots & & & & & & & & \\ \lambda_i \overset{1}{\tilde{\gamma}}^i &= 0, & \lambda_i \overset{2}{\tilde{\gamma}}^i &= 0, & \dots & \lambda_i \overset{r}{\tilde{\gamma}}^i &= 1, & \dots, & \lambda_i \overset{n}{\tilde{\gamma}}^i &= 0, \\ & \vdots & & & & & & & & \\ \lambda_i \overset{1}{\tilde{\gamma}}^i &= 0, & \lambda_i \overset{2}{\tilde{\gamma}}^i &= 0, & \dots & \dots & \dots, & \lambda_i \overset{n}{\tilde{\gamma}}^i &= 1, & [3]. \end{aligned}$$

We arrive at the conformal Frenet formulae

$$(3.2) \quad \begin{aligned} \frac{D\overset{1}{\gamma}^i}{d\kappa} &= \overset{2}{\gamma}^i, & \frac{D\lambda_i}{d\kappa} &= \lambda_i, \\ \frac{D\overset{2}{\gamma}^i}{d\kappa} &= -k_1^* \overset{1}{\gamma}^i - k^{**} \overset{2}{\gamma}^i - \bar{k}_2^* \overset{3}{\gamma}^i, & \frac{D\lambda_i}{d\kappa} &= -\bar{k}_1^* \lambda_i + k_2^* \lambda_i + k_3^* \lambda_i, \\ \frac{D\overset{3}{\gamma}^i}{d\kappa} &= \bar{k}_2^* \overset{1}{\gamma}^i + \bar{k}^* \overset{4}{\gamma}^i, & \frac{D\lambda_i}{d\kappa} &= \bar{k}_2^* \lambda_i + k_4^* \lambda_i, \\ \frac{D\overset{r}{\gamma}^i}{d\kappa} &= -k_{r-1}^* \overset{r-1}{\gamma}^i + \bar{k}_r^* \overset{r+1}{\gamma}^i, & \frac{D\lambda_i}{d\kappa} &= -\bar{k}_{r-1}^* \lambda_i + k_{r+1}^* \lambda_i, \\ \frac{D\overset{n}{\gamma}^i}{d\kappa} &= -k_{n-1}^* \overset{n-1}{\gamma}^i, & \frac{D\lambda_i}{d\kappa} &= -\bar{k}_{n-1}^* \lambda_i, \end{aligned}$$

(4 ≤ r ≤ n-1),

where

$$\frac{dk^*}{d\kappa} = k_1^* - \bar{k}_1^*, \quad k_2^* \bar{k}_2^* = k^* (\bar{k}_1^* + k_1^*) - k^{**} - \lambda_i^{(2)} \overset{1}{\gamma}^i, \quad \lambda_i^{(2)} = \lambda_i \overset{2}{\gamma}^i$$

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$$\bar{k}^{\epsilon} k_{\epsilon}^{\delta} = \sum_{s=3}^r (-1)^{s-1} \lambda_s^{(\epsilon)} \eta^{\delta(s)}, \quad (3 \leq r \leq n-1), \quad [3].$$

In these equations the n vectors $\overset{1}{\eta}^{\epsilon}, \overset{2}{\eta}^{\epsilon}, \dots, \overset{n}{\eta}^{\epsilon}$ form the associate conformal ennuple at a point of the curve.

Thus, making use of the quantity :

$$J^*(\kappa) = J \frac{d^2 J}{d\kappa^2} - \frac{3}{2} \left(\frac{dJ}{d\kappa} \right)^2 - \frac{1}{2} \left\{ Q(\kappa) - 2J \frac{dQ}{d\kappa} \right\} / J^2,$$

$$Q(\kappa) = \frac{1}{p-6} \left\{ J^{p-1} (G_{\epsilon j}(\kappa) \eta^{\epsilon} \eta^j + (2p-3)JK \right\},$$

the existence theorem in the conformal theory of a curve can be treated by the same method as that A. Fialkow [6].

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